

Cosmic Strings and Closed time-like curves in teleparallel gravity

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Abstract

Closed Time-like curves (CTC) in Cosmic strings in teleparallel T_4 gravity are forbidden. This result shown here in T_4 was shown by Soleng (Phys.Rev.D49 (1994)1124) also to be valid in Einstein-Cartan (EC) gravity. Here we show that in T_4 to allow for CTC we are also led to a lower bound on the angular momentum of the cosmic string. This result is obtained by matching the interior T_4 solution to a General Relativity (GR) vacuum solution. One of the main differences of the present report and the one by Soleng is that here the interior symmetric solution does not have necessary polarized spins but only Cartan torsion in the spirit of teleparallelism. Torsion flux is computed and it is shown that the center of cylinder singularity corresponds to a $2+1$ spacetime rotating point particle in T_4 . Therefore the possibility of building time machines seems to be strongly constrained than in the case of EC gravity.

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Recently H.Soleng and Letelier [1, 2, 3, 4] have investigated spinning cosmic string solutions in EC gravity. Letelier solution is easily shown to be trivial in Einstein teleparallelism [5]. Soleng solution is based on the construction of a spin polarized cylinder which allows us to build a spinning cosmic string with torsion where only one component of torsion survives and where the spinning particles are direct along the infinite axis of the cylinder in Einstein-Cartan theory. Soleng showed based on this model that close to the cosmic spinning string (or near a point particle equations. In this report we show that in T_4 CTC curves would imply on a lower bound constraint on the cosmic string angular momentum in order to allow for CTC and consequently to build time-machines. Teleparallel solutions have recently attracted attention since the works of J.G.Pereira and his group [6, 7] who show that it is possible to explain old problems in physics such as the tetrad complex of gravitational energy by making use of teleparallelism [6] as well as build some new teleparallel Schwarzschild and Kerr solutions [7] and discussing the Lense-Thirring effect in teleparallel geometry. More recently yet Garcia de Andrade has shown [8, 9] that is possible to explain another problems in cosmology such as the global rotation of the universe by using Gödel teleparallel solution in a very natural way. Besides as shown very recently it is possible to obtain a system of a cosmic string inside a black hole as a T_4 solution. Let us begin by considering the spacetime metric with cylindrical symmetry considered by Soleng as given by the one-form basis

$$\omega^0 = dt + M d\phi \quad (1)$$

$$\omega^1 = dr \quad (2)$$

$$\omega^2 = \rho d\phi \quad (3)$$

$$\omega^3 = dz \quad (4)$$

In terms of Cartan exterior differential forms the Soleng metric can be expressed as

$$ds^2 = (\omega^0)^2 - (\omega^1)^2 - (\omega^2)^2 - (\omega^3)^2 \quad (5)$$

Now from the Cartan structure equations

$$Q^i = d\omega^i + \omega_j^i \wedge \omega^j \quad (6)$$

where ω_j^i is the connection one-form and

$$R_j^i = R_{jkm}^i(\Gamma)\omega^k \wedge \omega^m = d\omega_j^i + \omega_k^i \wedge \omega_j^k \quad (7)$$

Here \wedge is the exterior product of forms symbol and $R_{jkl}^i(\Gamma)$ are the components of the Riemann-Cartan geometry curvature tensor and R_j^i is the curvature 2-form. Rewriting the metric (1) in the differential forms language one obtains

$$ds^2 = \eta_{ij}\omega^i\omega^j \quad (8)$$

where $\eta_{ij} = \text{diag}(+1, -1, -1, -1)$ is the tetrad Minkowski metric. By making use of the teleparallel condition $R_{jkl}^i(\Gamma) = 0$ into the equation (7) we notice that the constraint

$$\omega_j^i = 0 \quad (9)$$

fulfills the teleparallel condition. Here we adopt this stronger teleparallel condition which also has been adopted by Letelier [5] in the construction of torsion loops in teleparallel spacetimes. By using the condition (9) into equation (6) one obtains the torsion 2-form in the form

$$Q^i = d\omega^i = T_{jk}^i\omega^j\wedge\omega^k \quad (10)$$

where T_{jk}^i are the components of the Cartan's torsion tensor. Applying this simple expression to the expressions for the basis one-forms of the cosmic spinning torsion string system above one obtains after a quick computation one obtains the components of the torsion tensor as

$$s_0 = T_{12}^0 = \frac{M_{,r}}{\rho} \quad (11)$$

$$s_1 = T_{12}^1 = \frac{\rho_{,r}}{\rho} \quad (12)$$

By assuming that both torsion components inside the cylinder are constants equations (11) and (12) become differential equations that can be solved immediatly to yield the interior solution

$$ds^2 = [dt + \frac{s_0}{s_1}e^{s_1 r}d\phi]^2 - [dr^2 + dz^2] - e^{2s_1 r}d\phi^2 \quad (13)$$

This metric has some very nice properties that now we wish to comment. First according to Paul Tod [10] classification of the conical metric

$$ds^2 = (dt + \alpha d\phi)^2 - \beta^2 r^2 d\phi^2 - (dz + \gamma d\phi)^2 \quad (14)$$

when the last term can be omitted but α is non-zero the metric would represent a point particle rotating in $2 + 1$ gravity. Therefore with this classification one could say by comparison with our interior solution (13) one notices that this solution could be interpreted as a singular behaviour at $r = 0$ where a point particle would rotate at the center of the cosmic string at $3 + 1$ gravity spacetime. Before making use of the matching conditions to discuss CTC conditions we could compute the torsion flux by making use of geometrical phases formula

$$\int_{\Sigma} Q^0 = \int \omega^0 = 2\pi \frac{s_0}{s_1} e^{s_1 r} \quad (15)$$

which shows that at the conic singularity $r = 0$ torsion flux is constant. CTC could be imagined as torsion loops around the cosmic strings [6]. Let us now point out that the interior solution to be physical should be matched to an exterior solution preferable with a GR vacuum of the type given by Soleng one-forms

$$\theta^0 = dt + \frac{J}{2\pi} d\phi \quad (16)$$

$$\theta^1 = dr \quad (17)$$

$$\theta^2 = \left(1 - \frac{\mu}{2\pi}\right)(r + r_0)d\phi \quad (18)$$

$$\theta^3 = dz \quad (19)$$

where J is the intrinsic angular momentum of the cosmic string and μ is the string mass. The Arkuszewski-Kopczynski-Ponomarev (AKP) [11] matching conditions of EC theory can be used here [2]

$$g_{22}|_+ = g_{22}|_- \quad (20)$$

$$g_{02,1}|_+ = g_{02,1}|_- - T_{012} \quad (21)$$

$$g_{22,1}|_+ = g_{22,1}|_- - 2T_{212} \quad (22)$$

where the plus sign refers to the exterior solution and the minus sign refers to the interior solution. Substitution of the respective above solutions into the matching conditions yields

$$e^{2s_1 r} = -\left(1 - \frac{\mu}{2\pi}\right)^2 (R + r_0)^2 \quad (23)$$

$$2\frac{s_0}{s_1}e^{s_1 R} = \frac{J}{\pi} - s_0 \quad (24)$$

where R is the radius of the cylinder. From these last two expressions allows us to determine the constant torsion components in terms of the cylinder radius. They also allows us to obtain a constraint between the angular momentum of the T_4 string and the ratio of the torsion components

$$\pi \frac{s_0}{s_1} = \frac{J}{[s_1 + 2e^{s_1 R}]} \quad (25)$$

Let us now from these expressions to examine the problem of construction of CTC in T_4 . We know from GR that the CTC are allowed with the following constraint

$$(1 - \frac{\mu}{2\pi})(r + r_0) < \frac{J}{2\pi} \quad (26)$$

which at the surface of the cylinder is

$$(1 - \frac{\mu}{2\pi})(R + r_0) < \frac{J}{2\pi} \quad (27)$$

To simplify from the matching conditions we reduce this expression to

$$\frac{s_0}{[1 + \sqrt{2\frac{s_0}{s_1}}]} < \frac{J}{2\pi} \quad (28)$$

This expression can be further simplified if we consider that torsion can be "isotropic" or that they have the same $s_0 = S_1$. To simplify from the matching conditions we reduce this expression to

$$\frac{s_0}{[1 + \sqrt{2}]} < \frac{J}{2\pi} \quad (29)$$

Nevertheless by a similar reasoning used by Soleng we can show that the assumption that CTC can be constructed leads us to a contradiction. Let us assume that $R \gg r_0$ and $\frac{\mu}{2\pi} \ll 1$ thus from expression (26) one obtains that

$$R < \frac{J}{2\pi} \quad (30)$$

Therefore if we consider that $R = \frac{h}{m}$ and $J = \mu \frac{h}{m}$ one may conclude that $\mu > 1$ which contradicts our assumption that $\mu 2\pi \ll 1$. Therefore one may conclude that there is no possibility of building CTC in T_4 as in the EC case.

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